

# Magnetospheric Physics

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# Outline

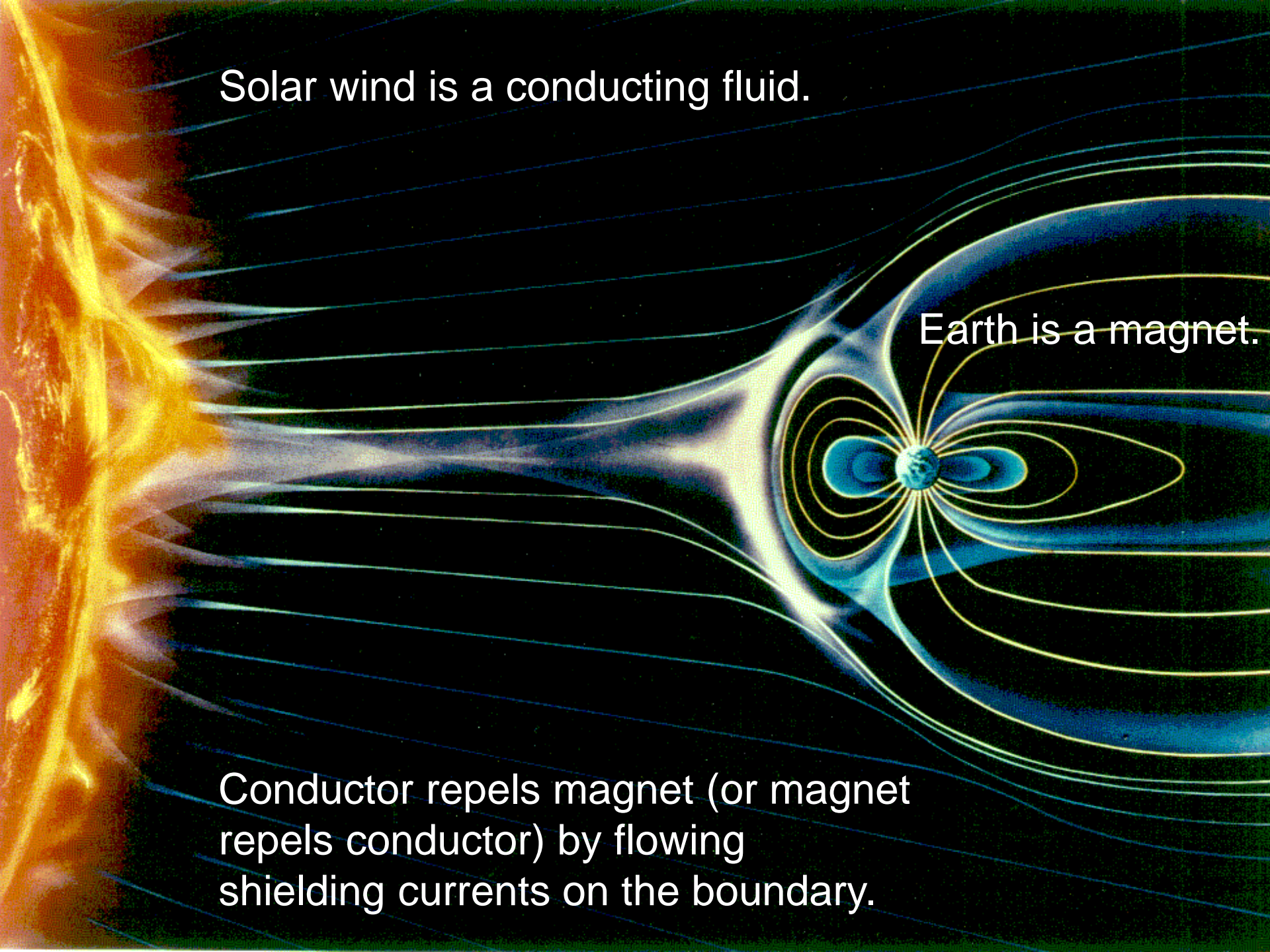
## Fundamentals of

1. large-scale structure and magnetic field topology of the magnetosphere
2. convection in the magnetosphere-ionosphere system

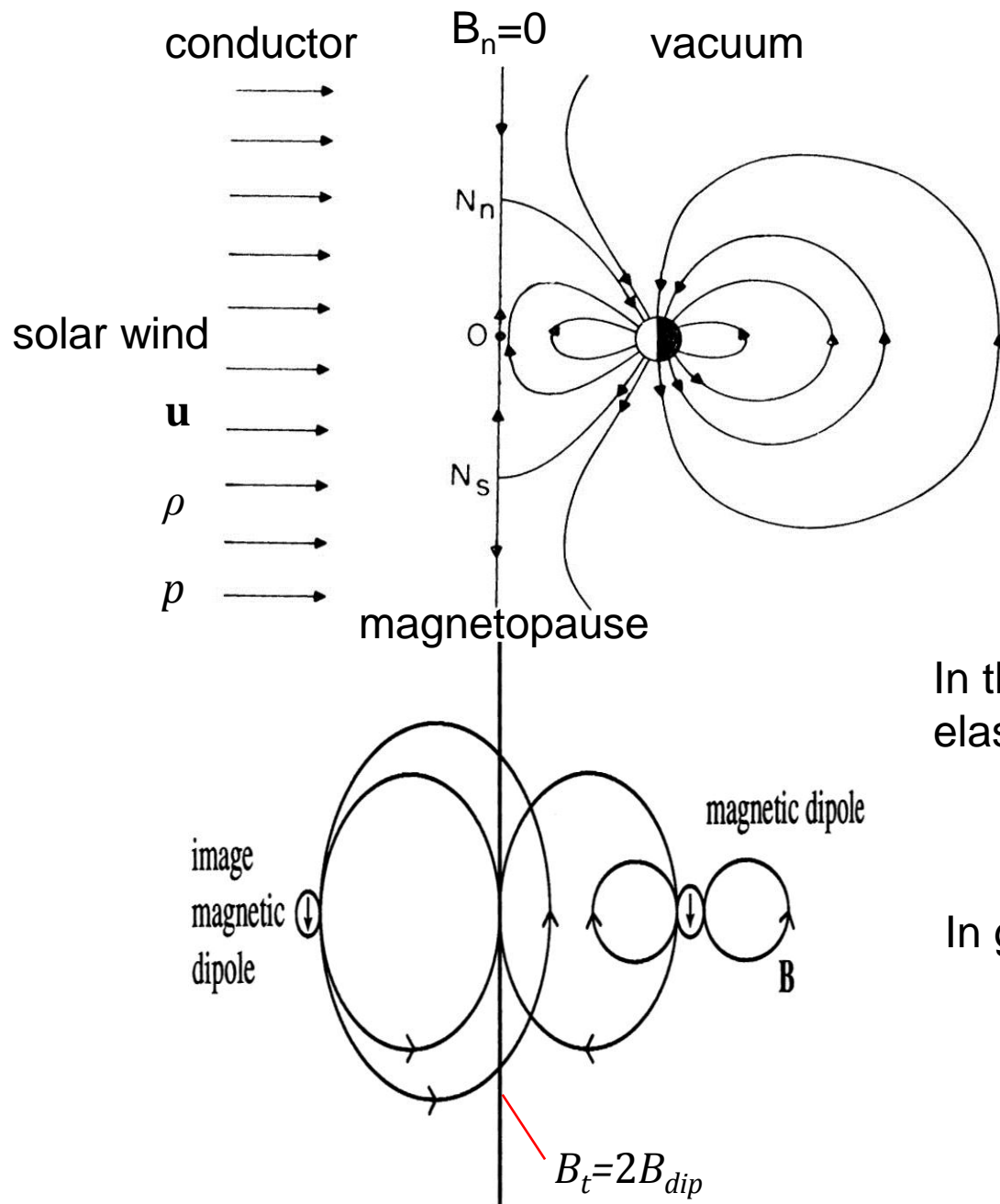
Solar wind is a conducting fluid.

Earth is a magnet.

Conductor repels magnet (or magnet repels conductor) by flowing shielding currents on the boundary.



# Vacuum model (Chapman & Ferraro)



Pressure balance

$$p_N = \frac{B_t^2}{2\mu_0}$$

$p_N$ : Newtonian pressure

$$p_N = \mathbf{n} \cdot (\rho \mathbf{u} \mathbf{u} + p \mathbf{I}) \cdot \mathbf{n}$$

$\mathbf{n}$ : boundary normal

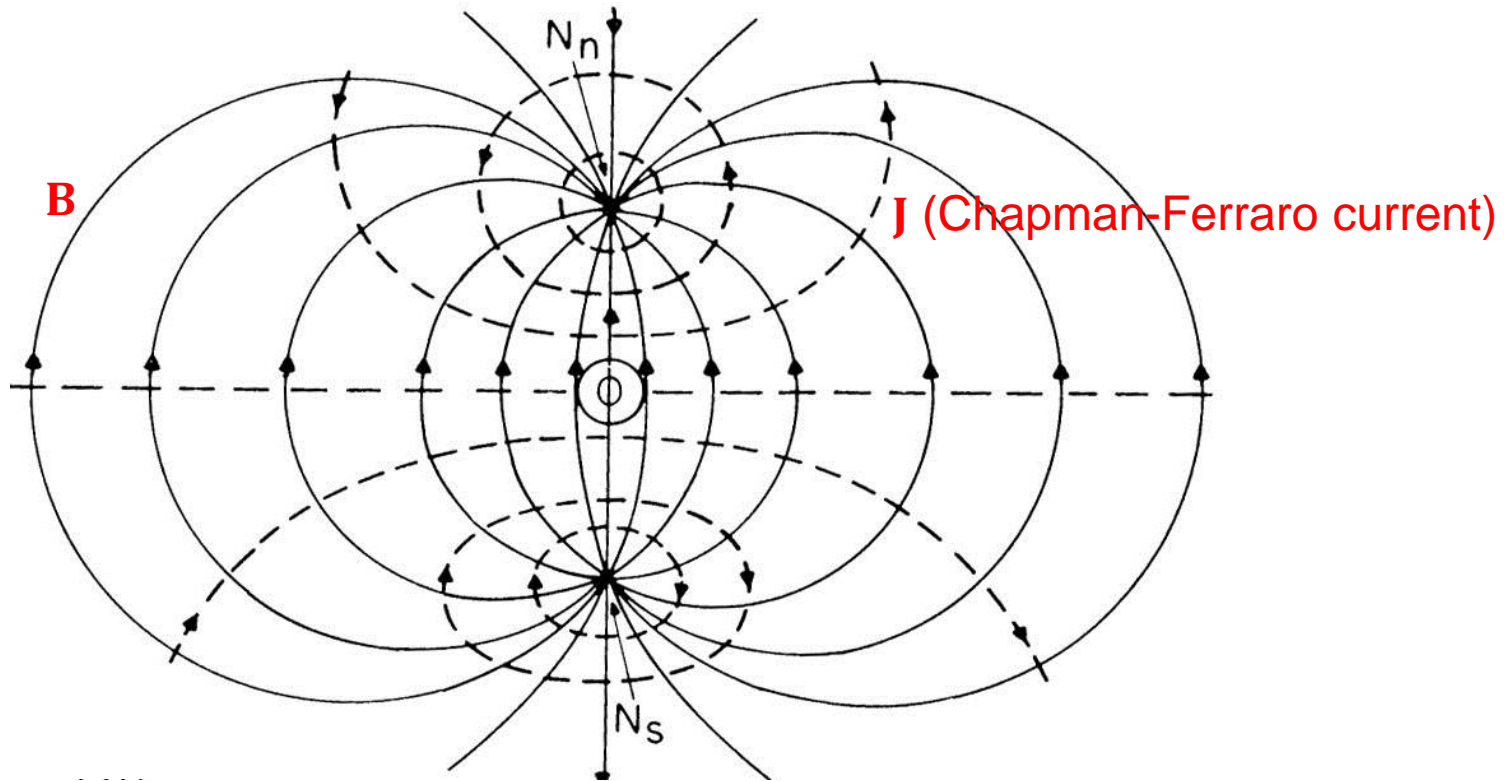
In the case of cold plasma beam that elastically reflects off the boundary

$$p_N = 2\rho u^2$$

In general

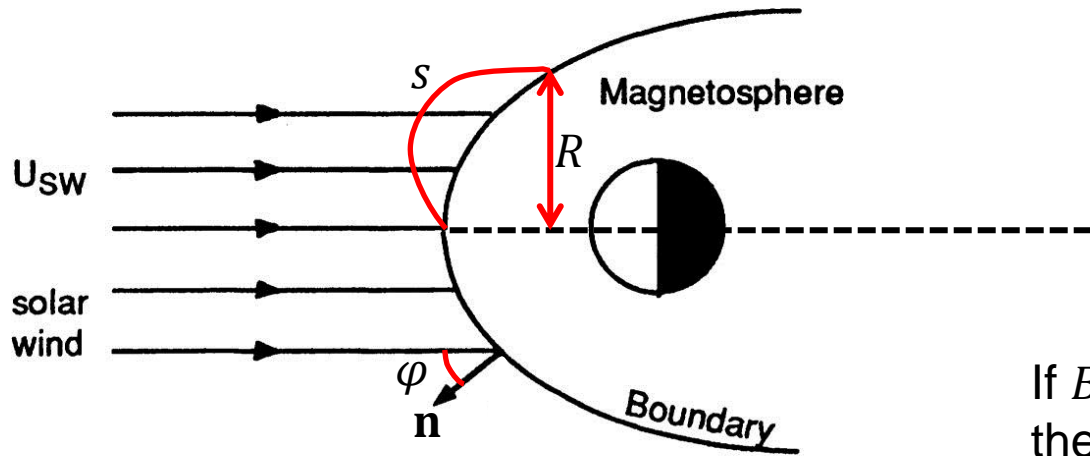
$$p_N = K\rho u^2 \quad (0 < K \leq 2)$$

# Chapman-Ferraro current



The role of Chapman-Ferraro current is to confine the Earth's magnetic field in a volume and at the same time to prevent the solar wind from penetrating into that volume.

## Realistic shape

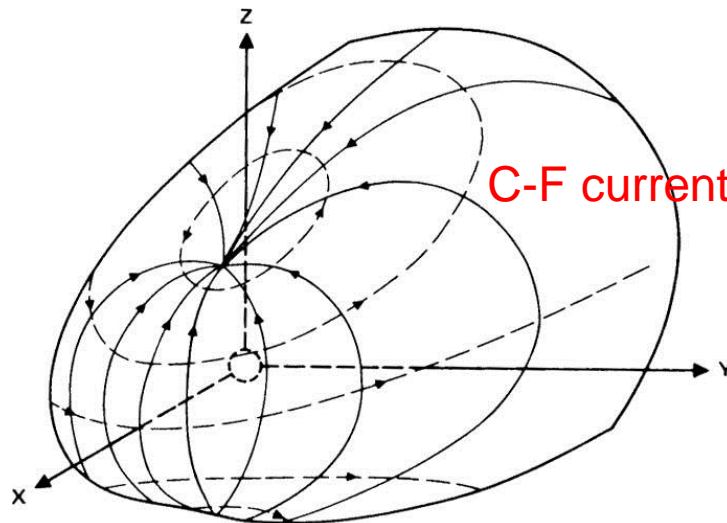


$$\left\{ \begin{array}{l} P_N = K\rho u^2 \cos^2 \varphi \\ p_N = \frac{B_t^2}{2\mu_0} \\ \frac{dR}{ds} = \cos \varphi \end{array} \right.$$

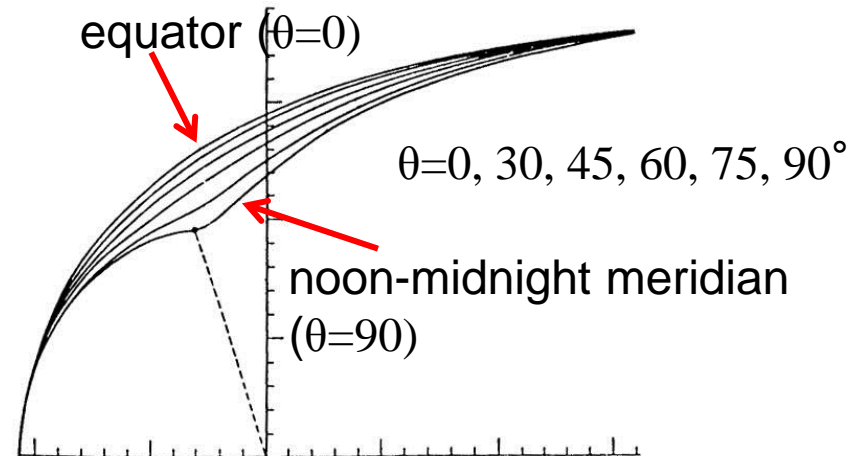
If  $B_t$  is given as a function of position, these equations determine the shape of the boundary.

$s$ : distance from the subsolar point along the boundary

## Examples of solutions



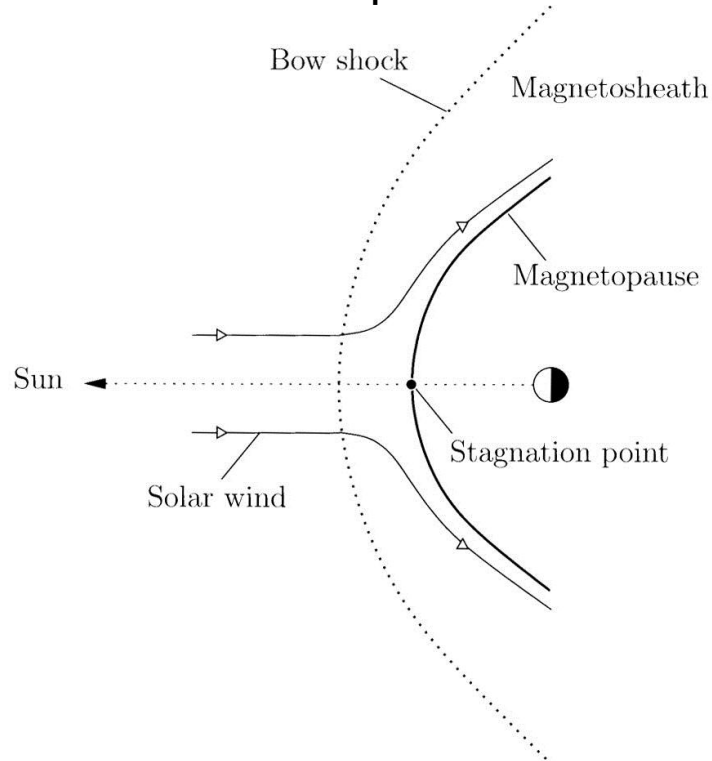
Midgley & Davis (1963)



Mead & Beard (1964)

$M_A \geq 10$  Magnetic field can be neglected.

$M \geq 10$  Thermal pressure can be neglected.



bow shock  $\rightarrow$  magnetosheath

$$\frac{u_2}{u_1} = \frac{\gamma - 1}{\gamma + 1} + \frac{2}{(\gamma + 1)M_1^2} \approx \frac{\gamma - 1}{\gamma + 1} = \frac{1}{4}$$

$$\frac{\rho_2}{\rho_1} = \frac{u_1}{u_2} = \frac{\gamma + 1}{\gamma - 1 + \frac{2}{M_1^2}} \approx \frac{\gamma + 1}{\gamma - 1} = 4$$

$$\frac{p_2}{p_1} = \frac{2\gamma M_1^2 - (\gamma - 1)}{\gamma + 1} \approx \frac{2\gamma M_1^2}{\gamma + 1} = \frac{2}{\gamma + 1} \frac{\rho_1 u_1^2}{p_1}$$

$$= \frac{2}{\gamma + 1} \frac{\rho_1 u_1^2}{p_1} = \frac{3}{4} \frac{\rho_1 u_1^2}{p_1} \rightarrow p_2 = \frac{3}{4} \rho_1 u_1^2$$

$$M_2^2 = \frac{2 + (\gamma - 1)M_1^2}{2\gamma M_1^2 - (\gamma - 1)} \approx \frac{\gamma - 1}{2\gamma} = \frac{1}{5}$$

$\rightarrow$  subsolar magnetopause

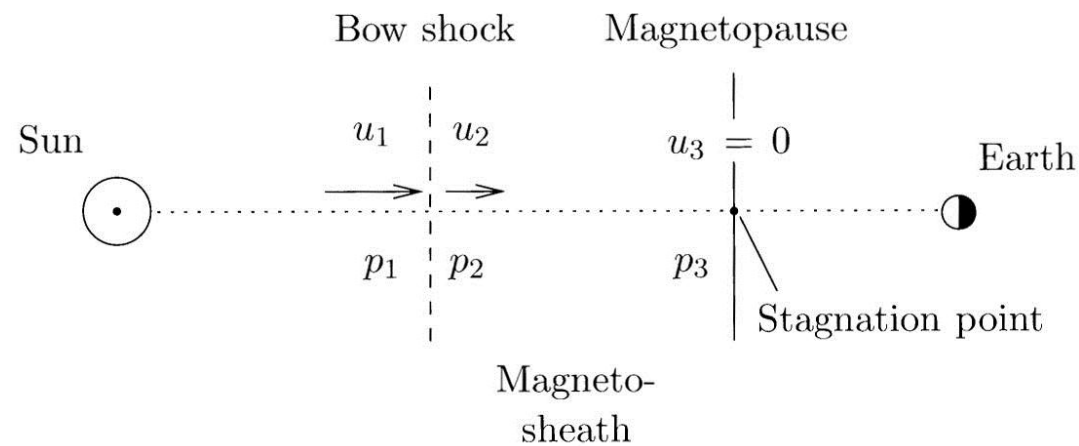
Bernoulli theorem

$$\frac{u_2^2}{2} + \frac{\gamma}{\gamma - 1} \frac{p_2}{\rho_2} = 0 + \frac{\gamma}{\gamma - 1} \frac{p_3}{\rho_3}$$

adiabatic  $p\rho^\gamma = \text{const.}$

$$\frac{p_3}{p_2} = \left(1 + \frac{\gamma - 1}{2} M_2^2\right)^{\frac{\gamma}{\gamma - 1}} = \left(\frac{16}{15}\right)^{\frac{5}{2}}$$

$$K = \frac{3}{4} \left(\frac{16}{15}\right)^{\frac{5}{2}} = 0.881$$

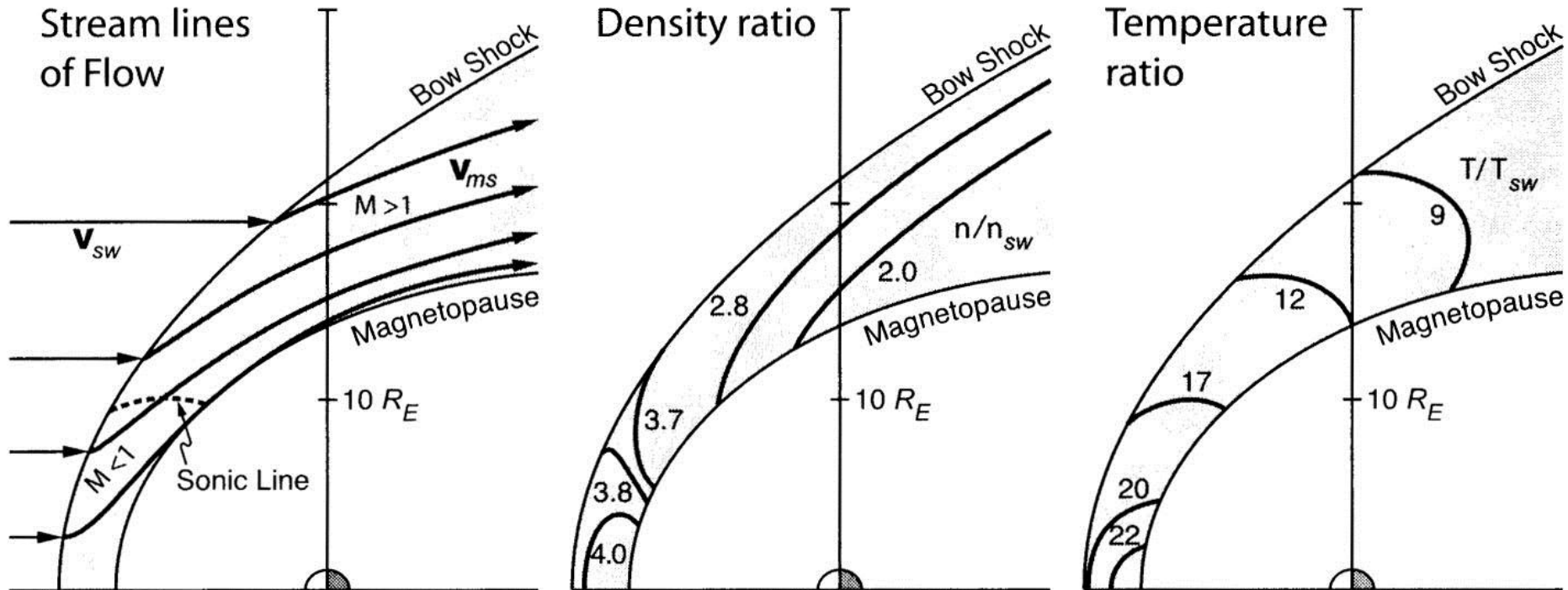


# Advanced vacuum model

When  $\varphi$  becomes large, the thermal pressure of the solar wind cannot be negligible.

$$P_N = K\rho u^2 \cos^2 \varphi + p_{sw}$$

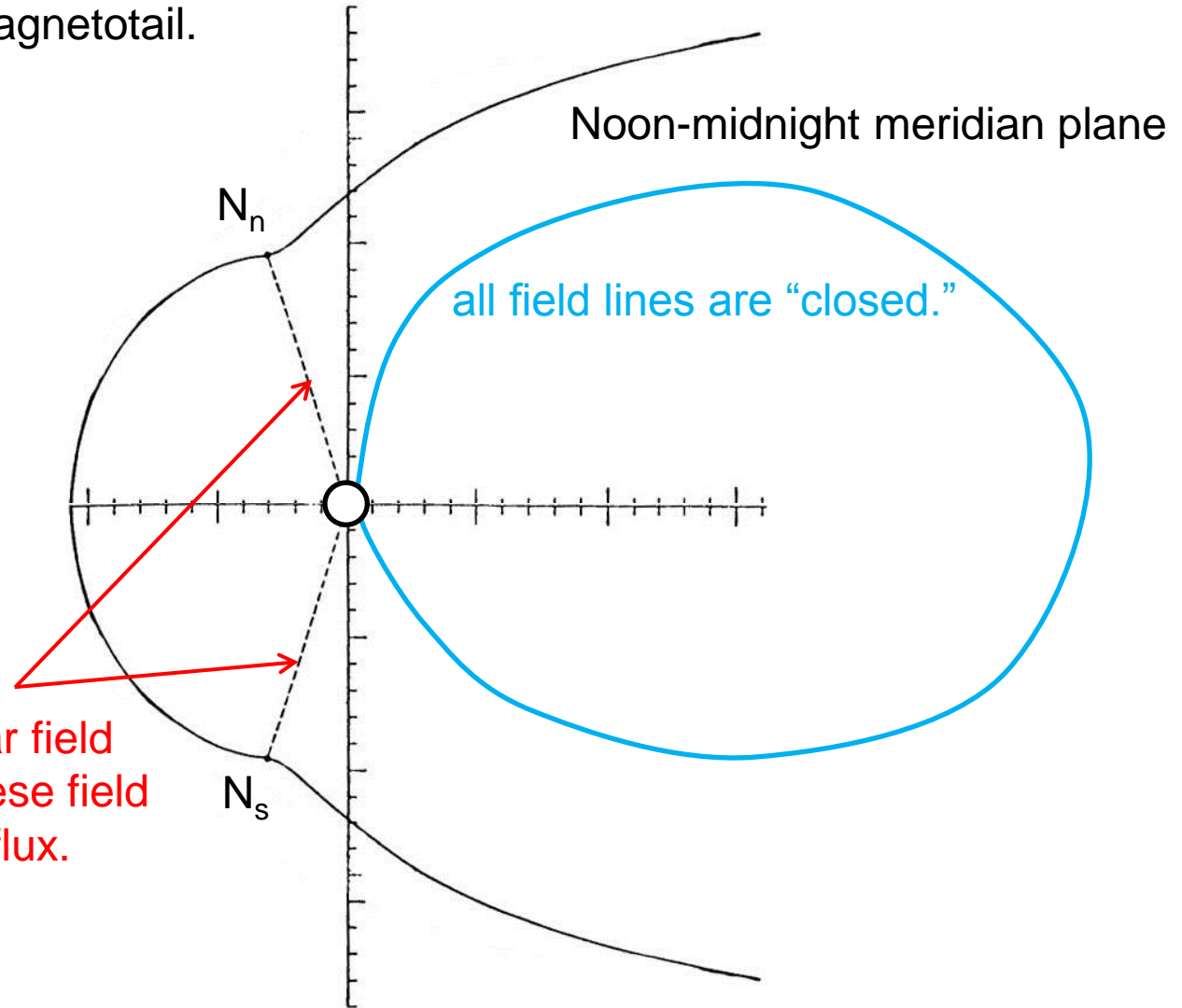
Numerical calculation for  $M=8$  (Spreiter et al., 1966)





# Magnetotail

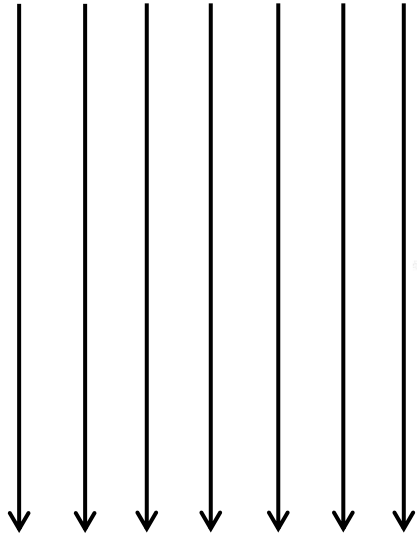
- Magnetotail may be defined as the region that conveys a finite magnetic flux from the Earth to infinity. Ionospheric projection of the magnetotail is called the polar cap.
- In the Chapman-Ferraro models (a vacuum magnetosphere confined to a volume), there is no magnetotail.



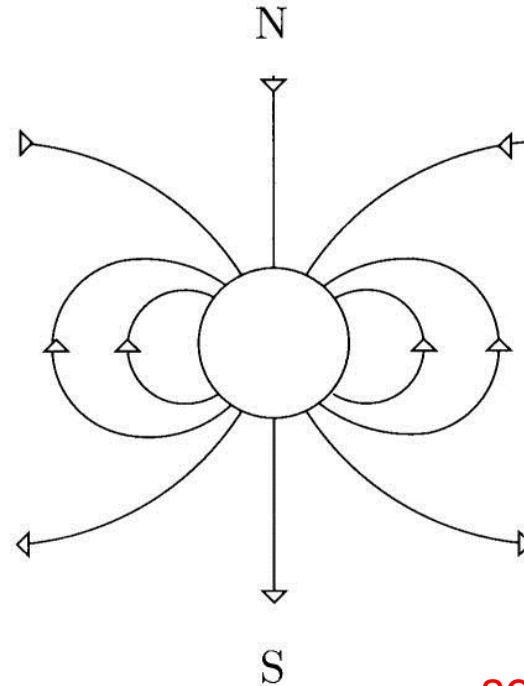
There are only two singular field lines that go to infinity. These field lines carry zero magnetic flux.

# Adding a magnetotail

Uniform southward field  
representing magnetotail



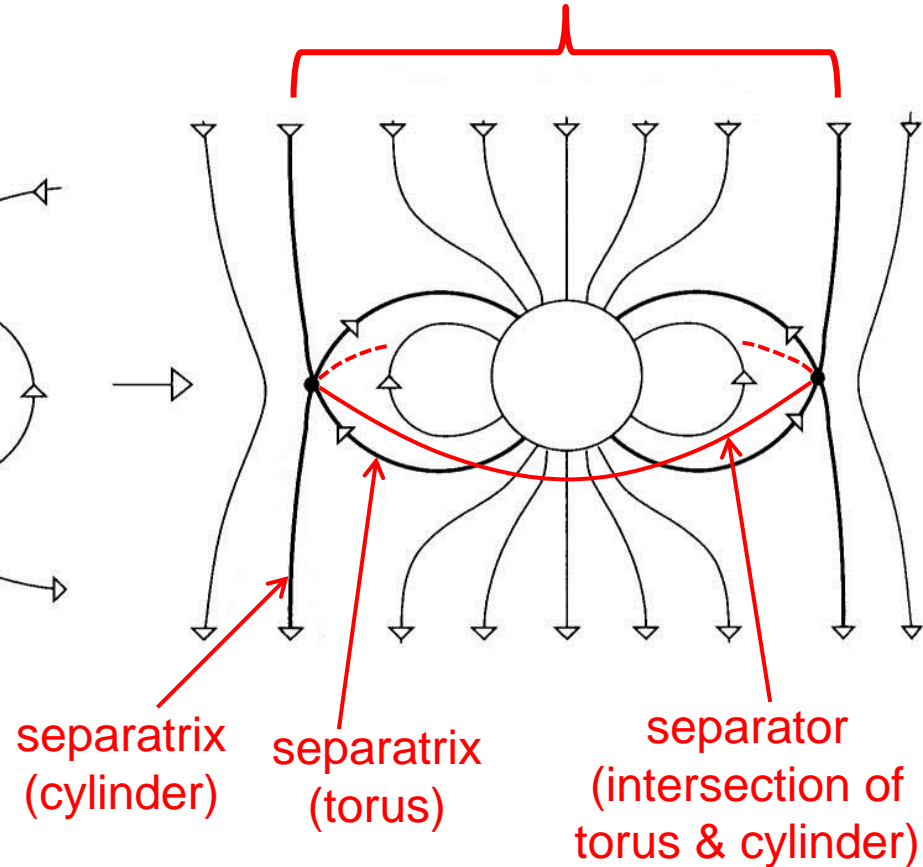
Earth's dipole field



+

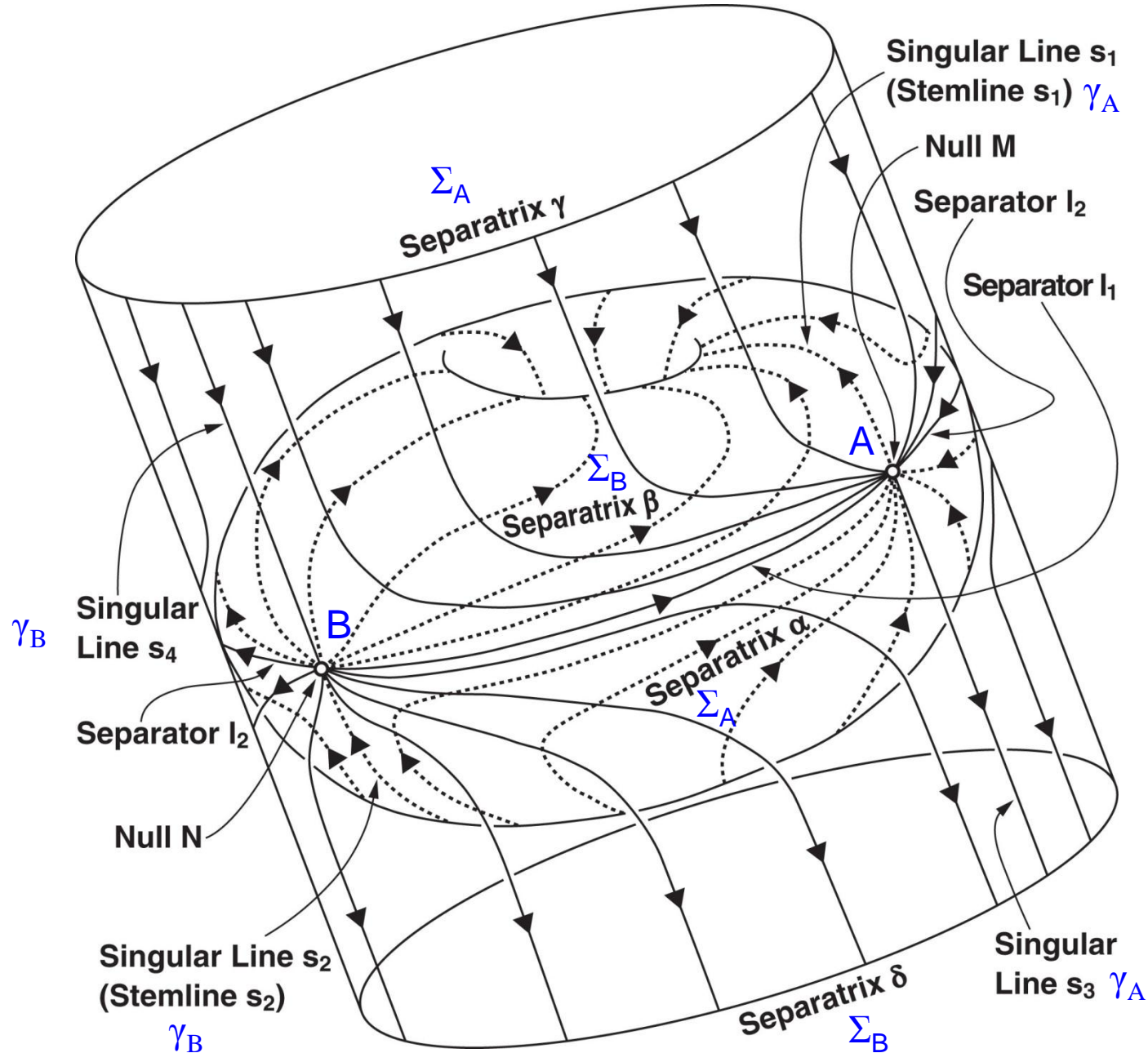


Magnetotail (open) field lines



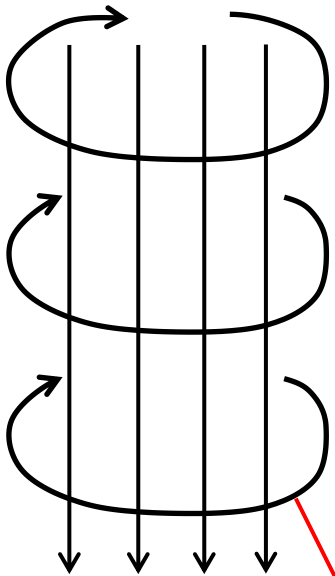
- Separatrix: magnetic surface that demarcates topologically different regions. It consists of field lines.
- Separator: intersection of separatrices. It is also a field line connecting two magnetic nulls.

# Superposition of dipole and nonparallel uniform fields



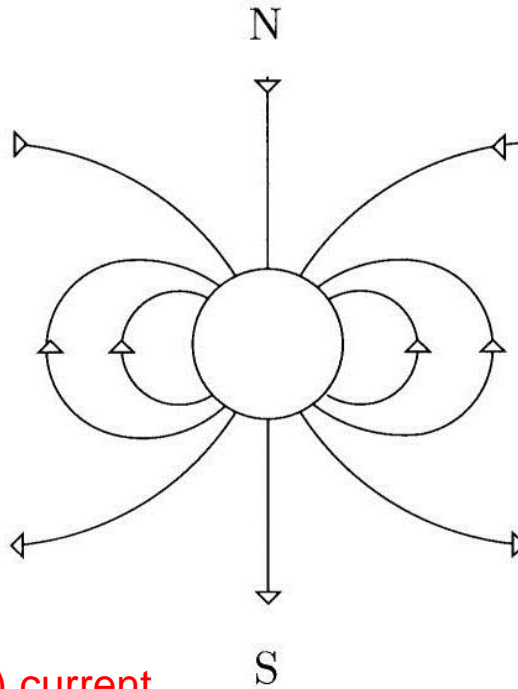
# Confined magnetotail (non-vacuum model)

confined tail field

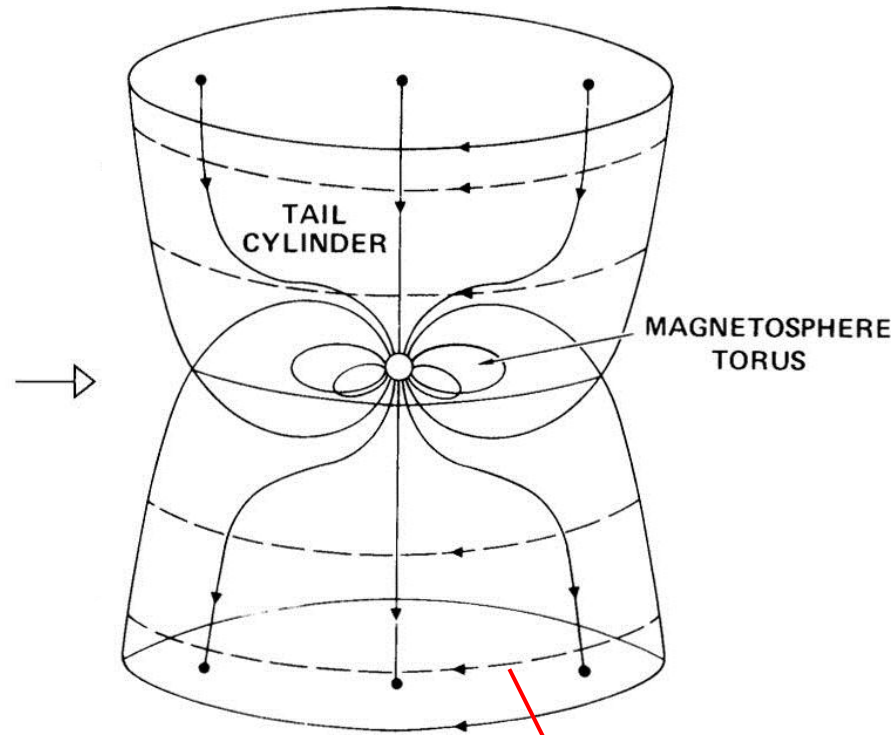


shielding (tail) current

Earth's dipole field



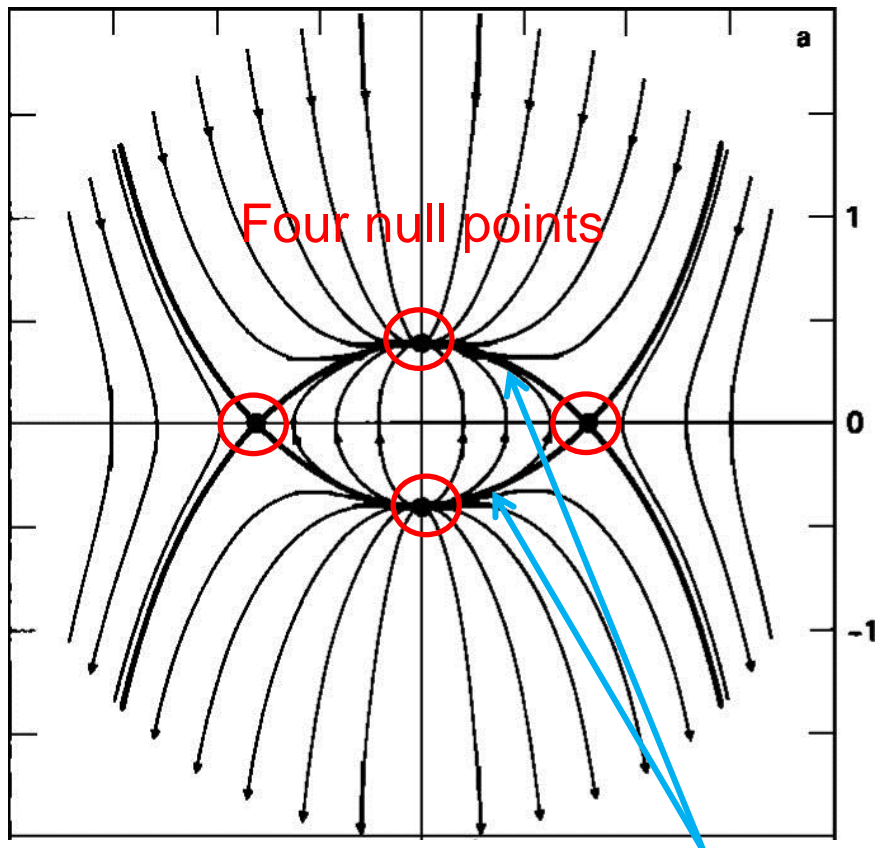
dipole + confined tail field



shielding (tail) current

# Confined tail flattened against an image plane

dipole + uniform field + image  
Crooker (1985)

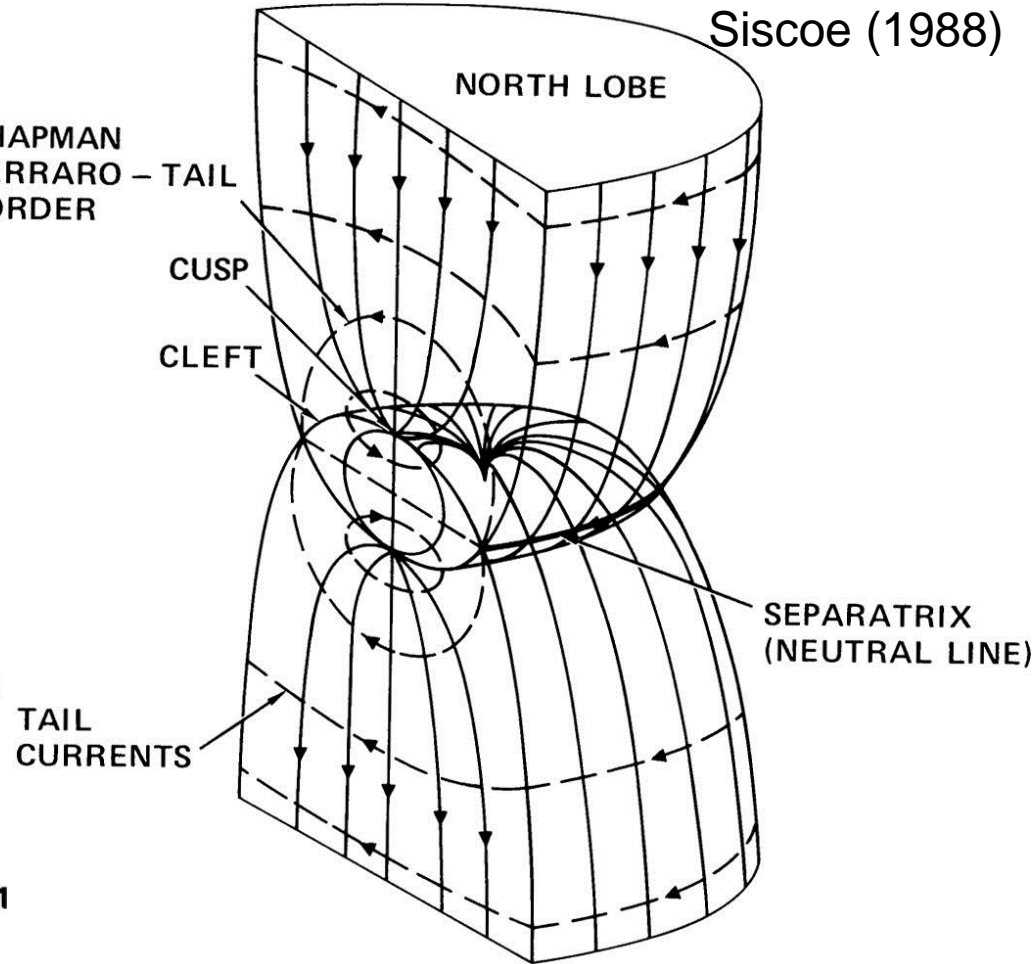


The separator splits into two.

dipole + confined tail + image

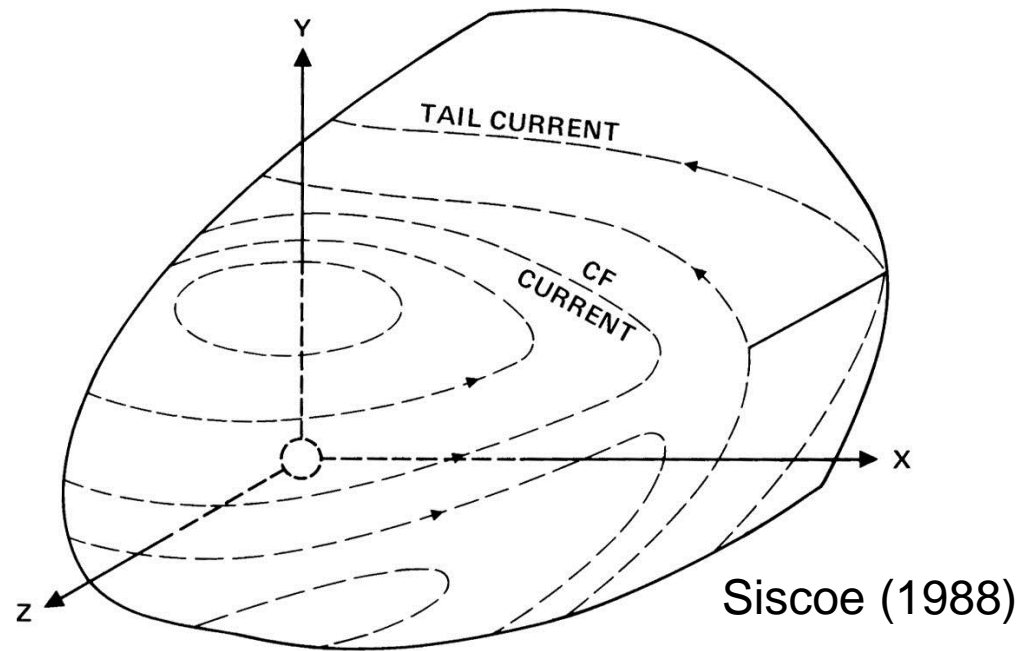
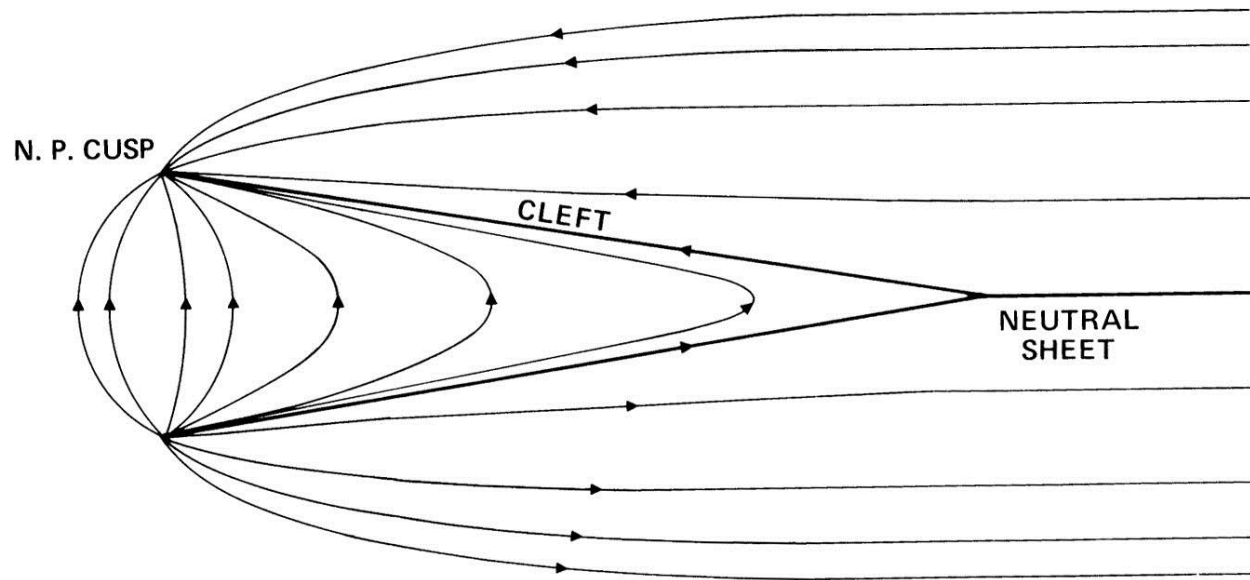
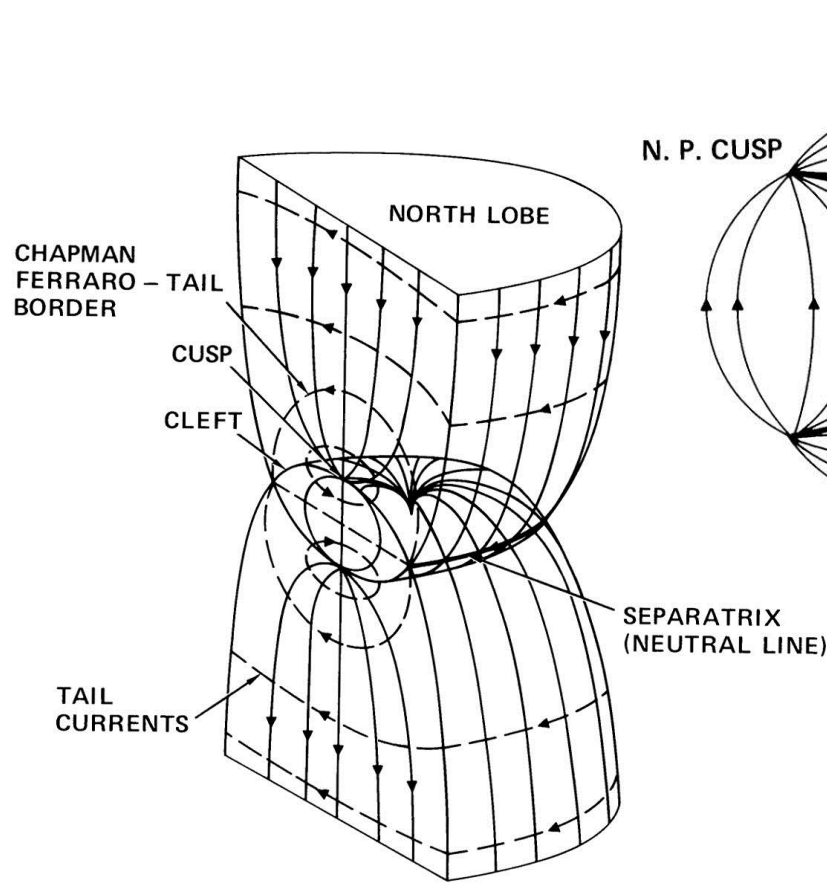
Siscoe (1988)

CHAPMAN  
FERRARO - TAIL  
BORDER



The C-F current is confined on the dayside. However, the C-F and tail currents are continuous on the magnetopause.

Final procedure is to mold the magnetosphere into a realistic shape.

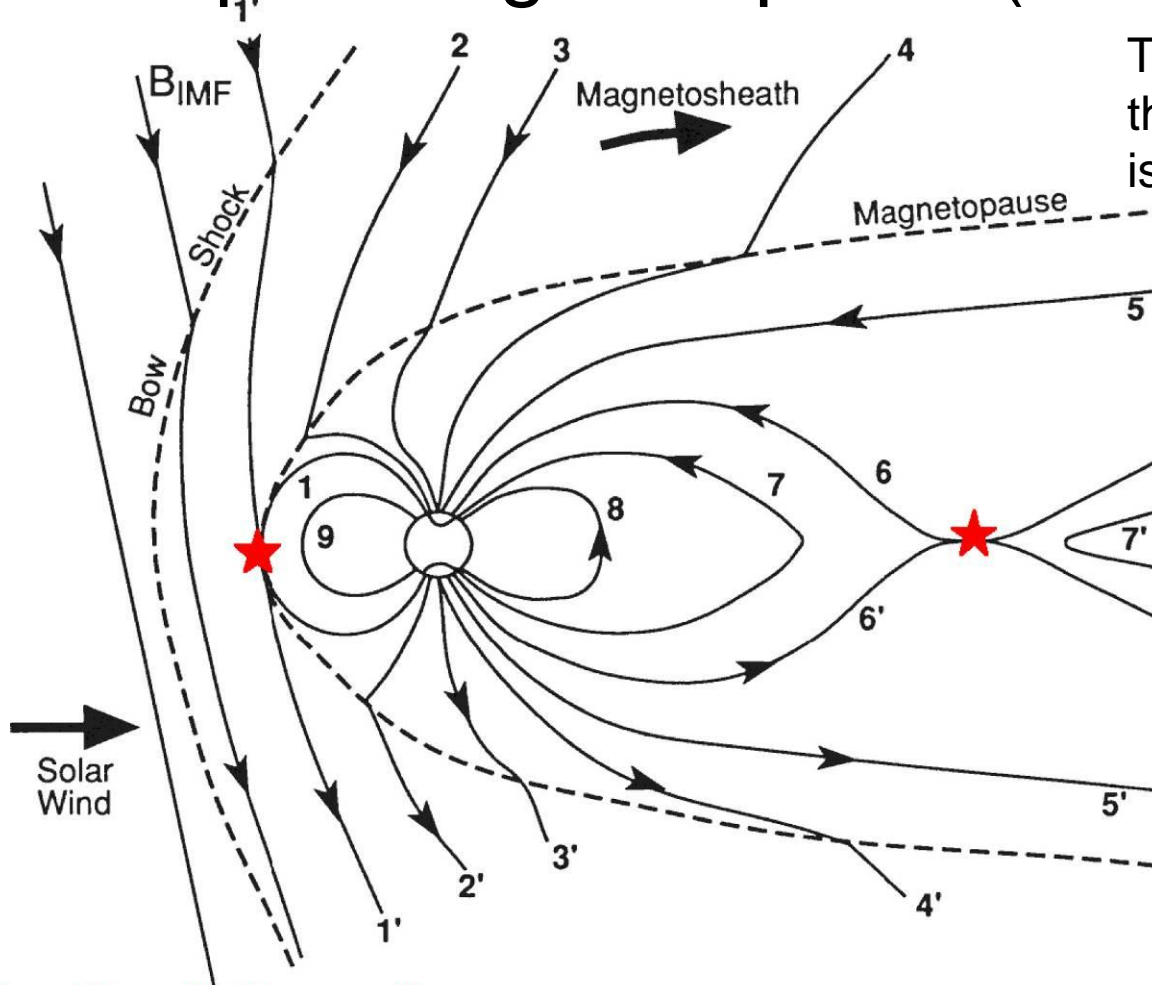


The magnetosphere is "closed" in the sense that the magnetopause is a tangential discontinuity.



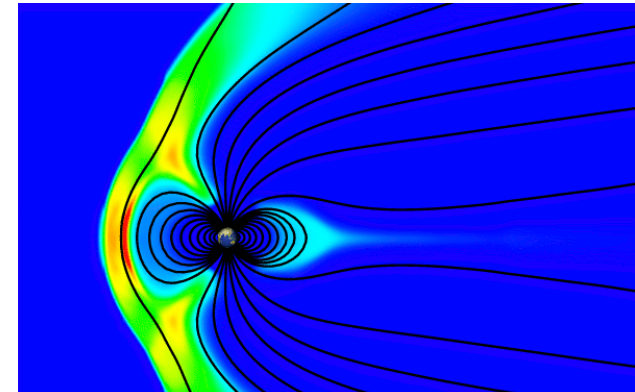
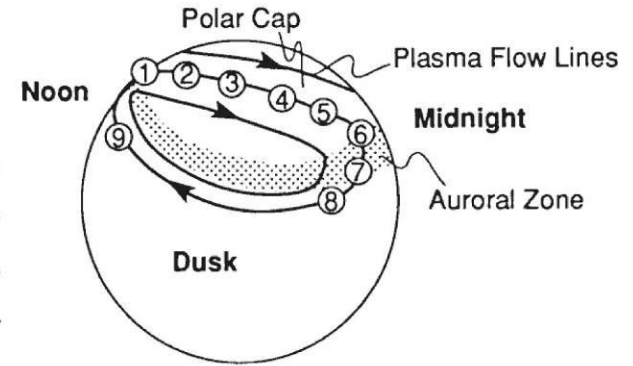
Connection to interplanetary magnetic field comes into play.

# Open magnetosphere (Dungey cycle)



The magnetosphere is “open” in the sense that the magnetopause is a rotational discontinuity.

## two-cell convection



★ --- Magnetic Reconnection

## Caveat

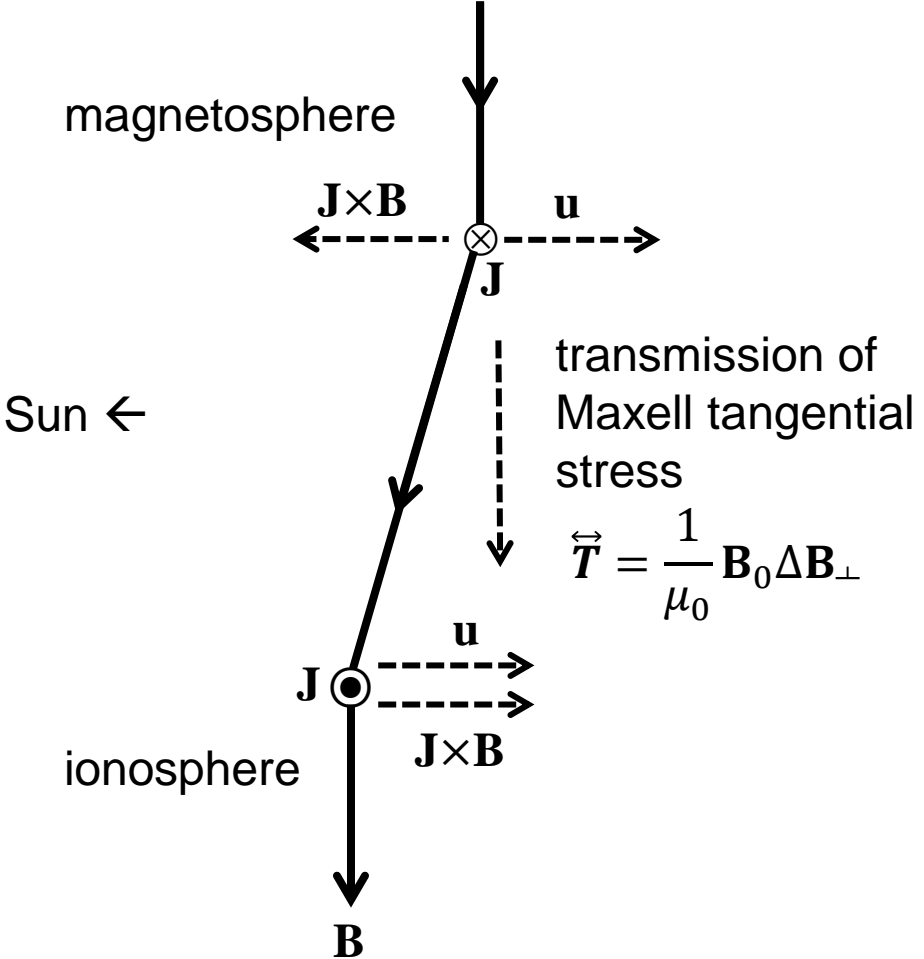
The figure gives an impression that the solar wind drags ionospheric plasmas to excite two-cell convection. However, this is not necessarily a correct view. **Ionosphere is very heavy.** There are complicated processes to establish a convection system (which we cannot review in detail here).

# Convection as a collaboration of magnetosphere & ionosphere (1)

## Momentum and energy

(a) ( $\mathbf{u}, \mathbf{B}$ ) view

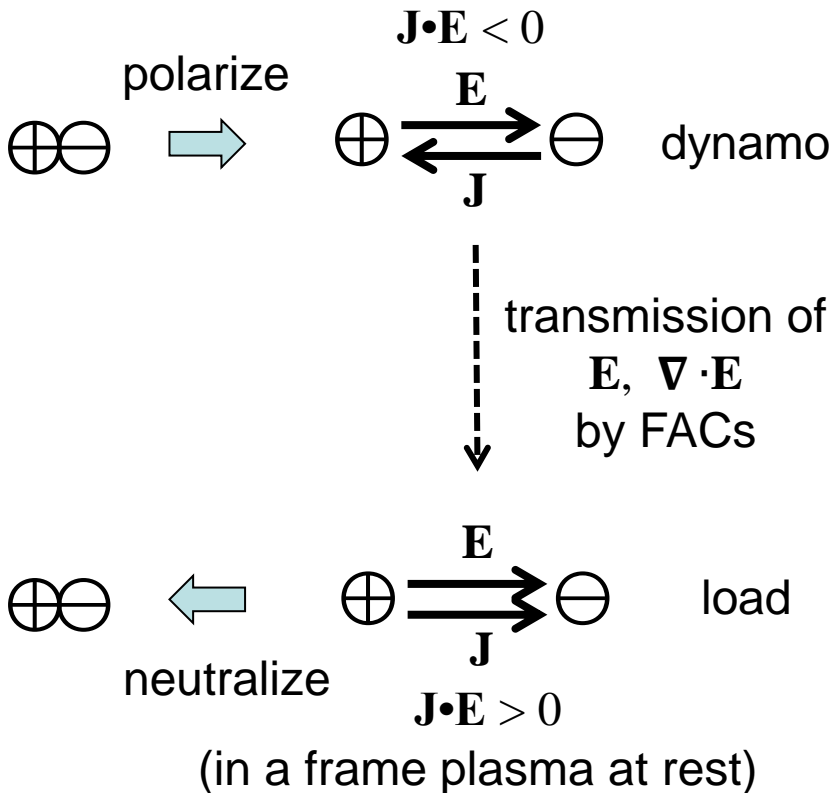
view from dusk



(b) ( $\mathbf{E}, \mathbf{J}$ ) view

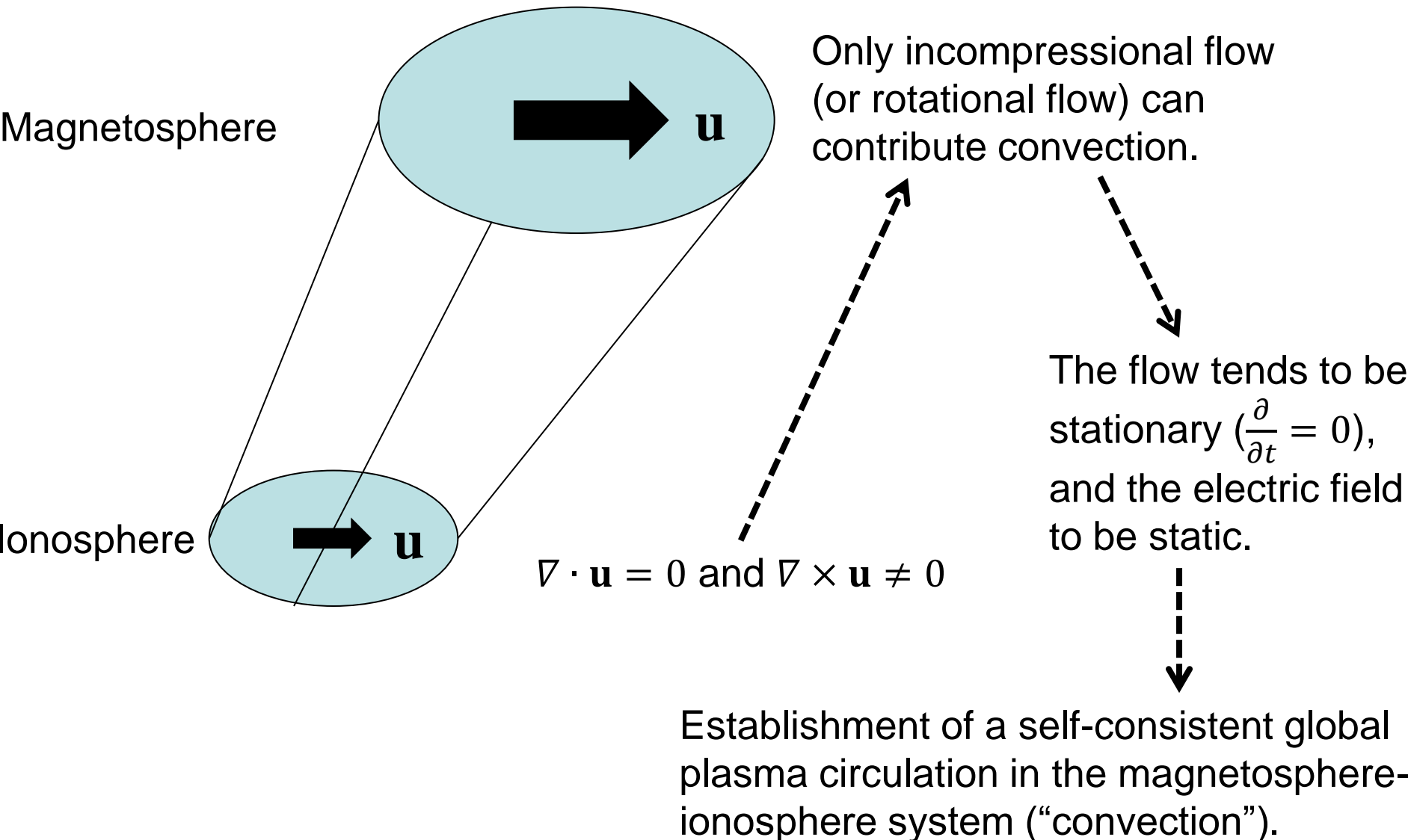
view from Sun

dawn  $\leftarrow$   $\rightarrow$  dusk

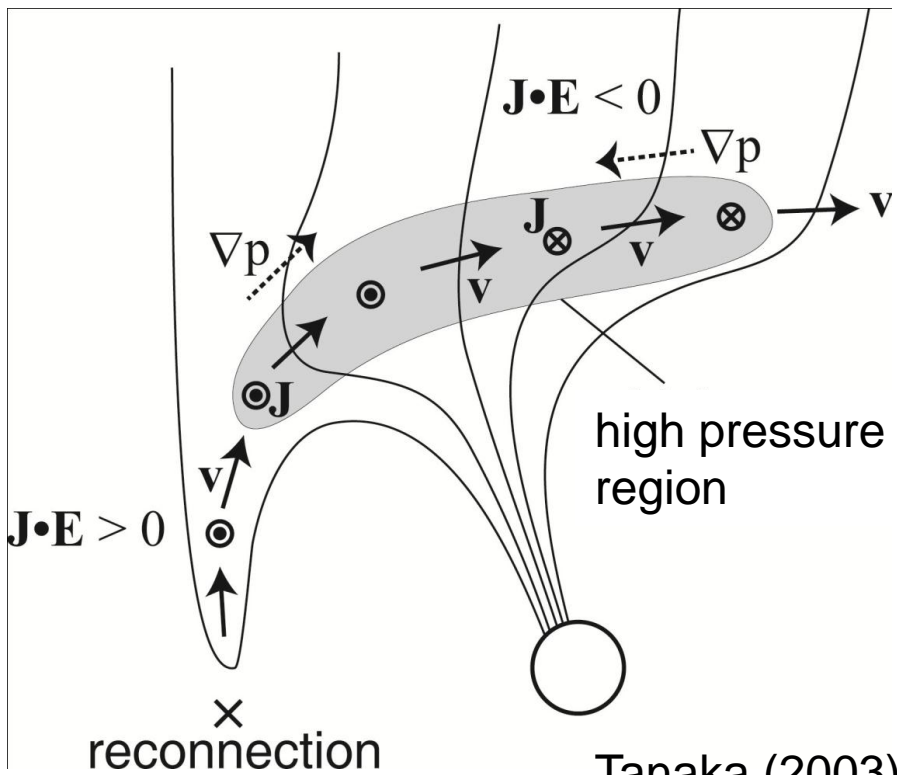
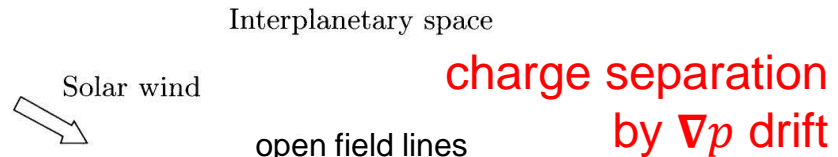
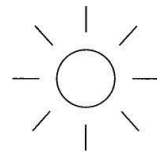




# Convection as a collaboration of magnetosphere & ionosphere (2)

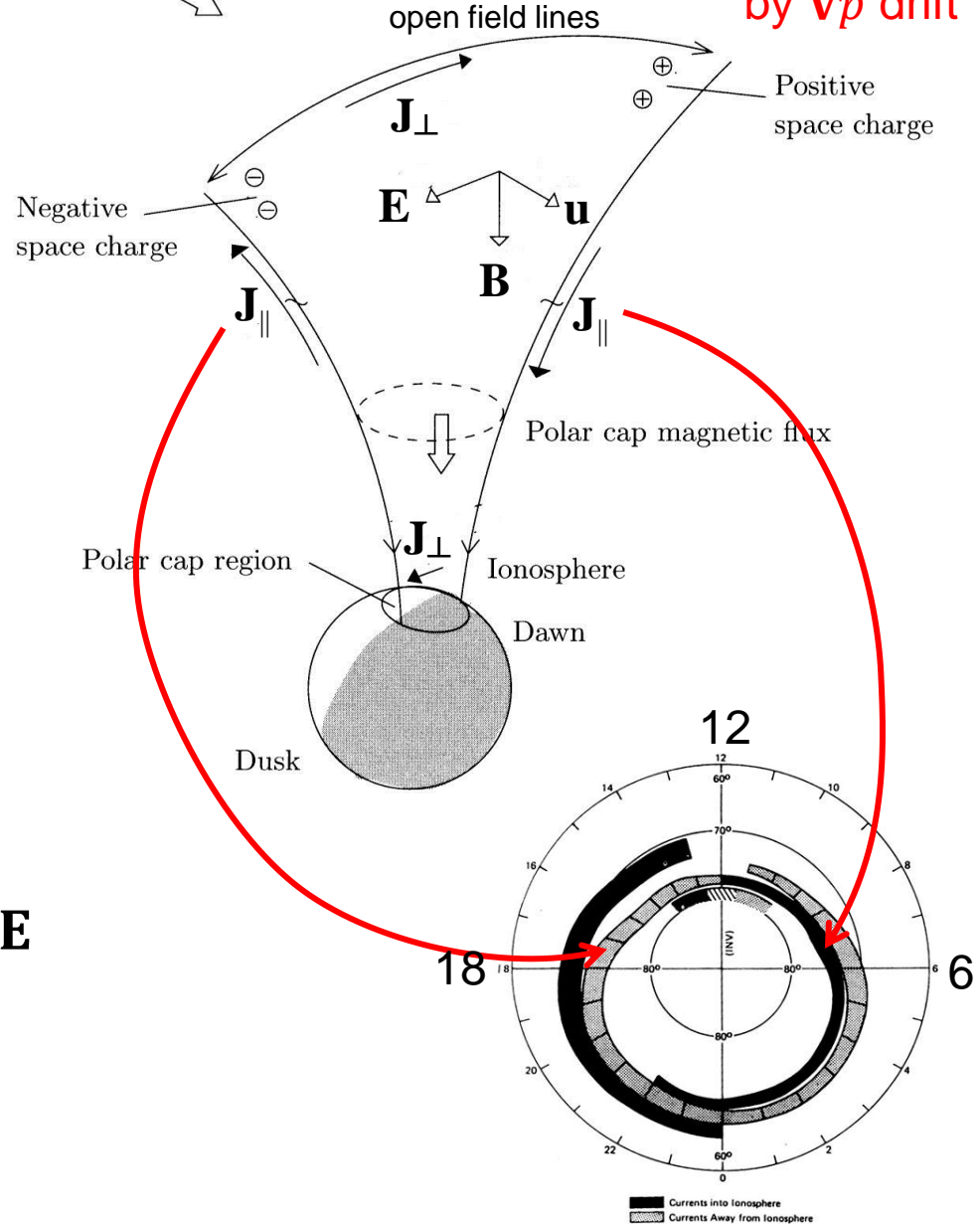


# Solar wind dynamo



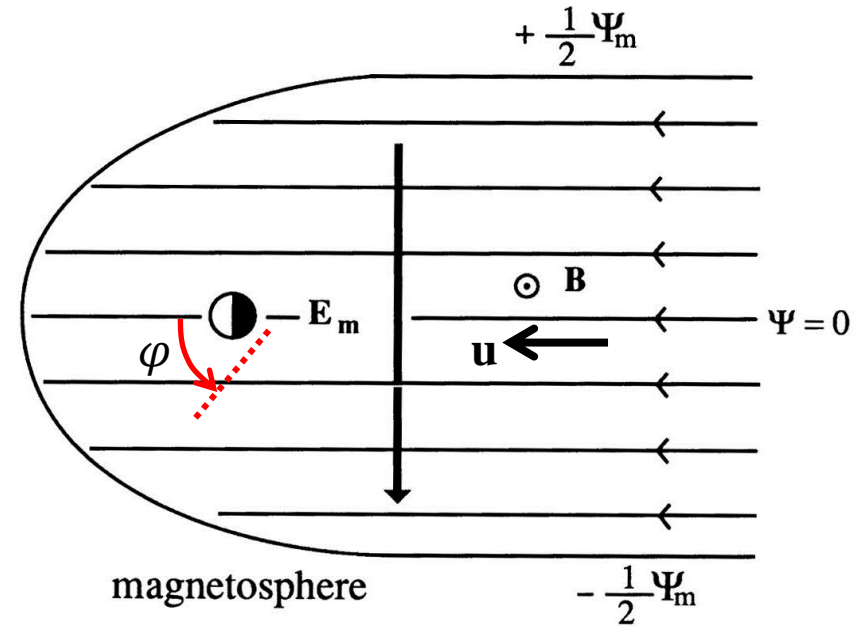
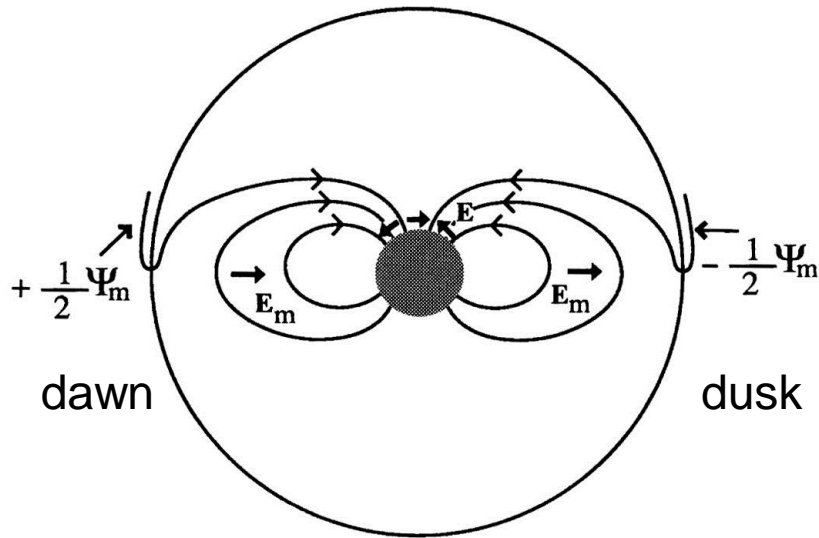
Tanaka (2003)

$$\mathbf{J} \times \mathbf{B} \approx \nabla p \quad \longrightarrow \quad \mathbf{v} \cdot \nabla p = \mathbf{J} \cdot \mathbf{E}$$



Currents into Ionosphere  
 Currents Away from Ionosphere

# Convection (Electric field) in the closed field line region



$$\mathbf{u} = \frac{\mathbf{E} \times \mathbf{B}}{B^2} \quad \mathbf{E} = -\nabla\Psi$$

(static electric field)

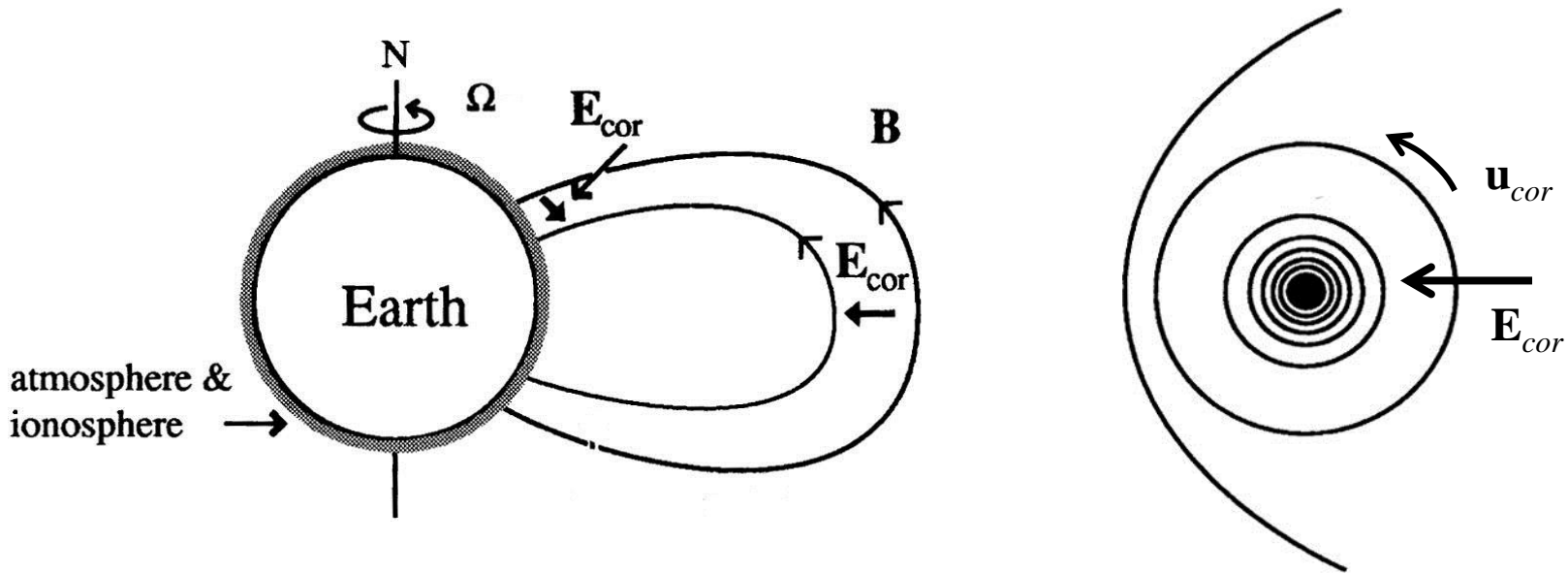
Uniform dawn-to-dusk electric field

$$\Psi_m = -E_m L a \sin \varphi \quad \mathbf{E}_m = -\nabla\Psi_m$$

$a$ : radius of the Earth

$L$ : geocentric distance measured by  $a$  where a field line crosses the equator

# Corotation electric field in the inner magnetosphere



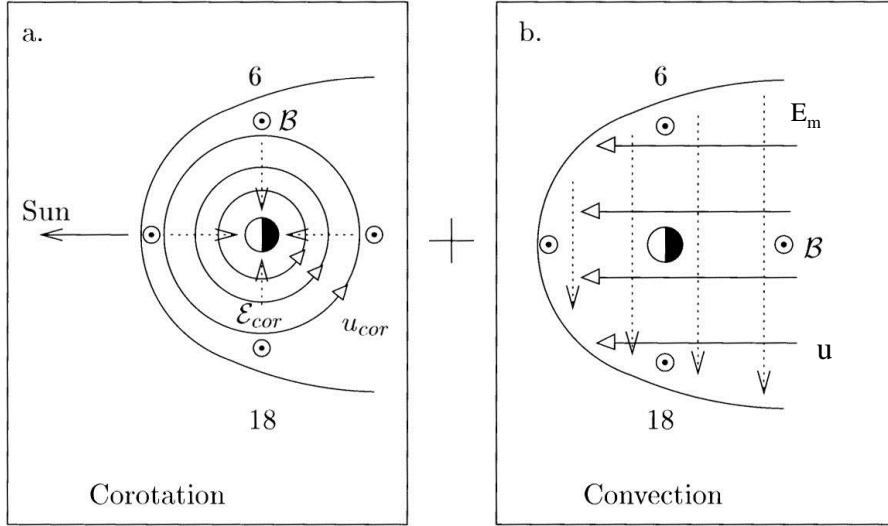
$$\mathbf{E}_{cor} = -\mathbf{u}_{cor} \times \mathbf{B} = -(\boldsymbol{\Omega} \times \mathbf{r}) \times \mathbf{B}$$

At the equator, assuming a dipole,

$$E_{cor} = -\frac{\Omega B_0 a}{L^2} \quad \text{radial component}$$

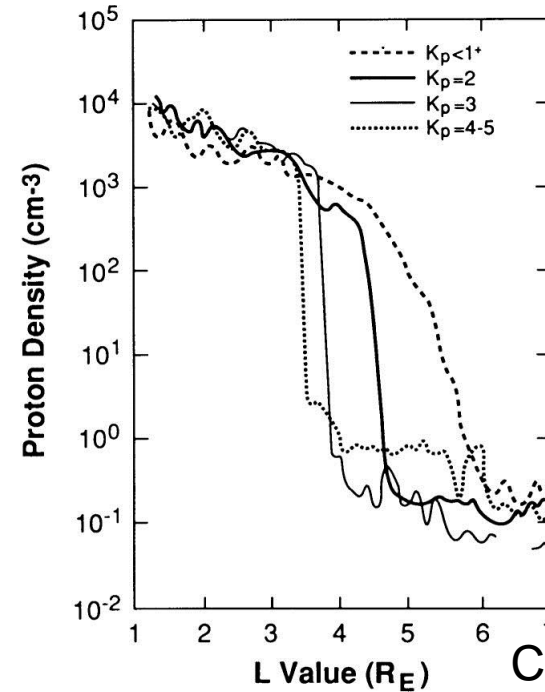
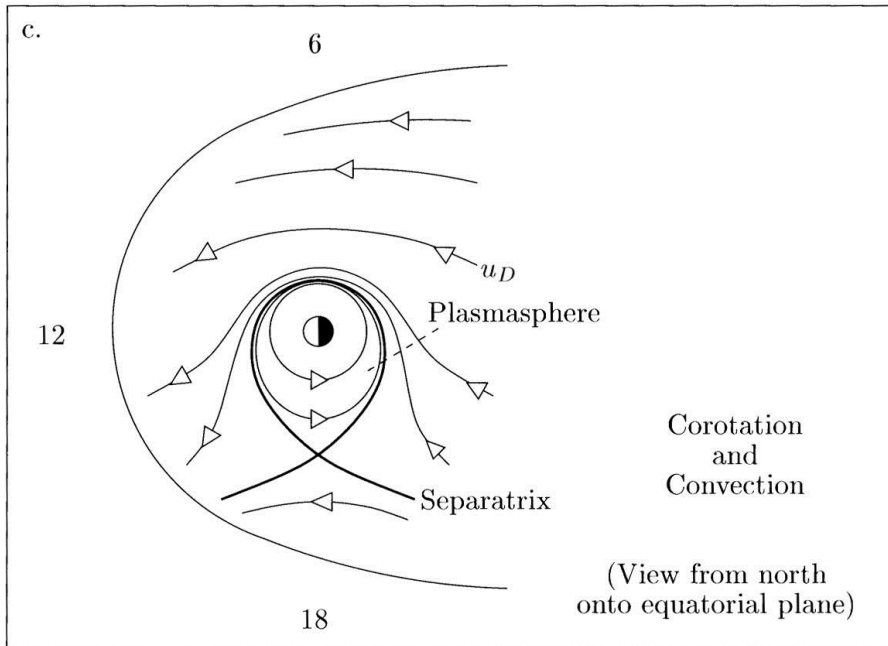
$$\Psi_{cor} = -\frac{a^2 \Omega B_0}{L}$$

# Total convection



$$\Psi = \Psi_m + \Psi_{cor} = -E_m L a \sin \phi - \frac{\Omega B_0 a^2}{L}$$

$$\mathbf{u} = \frac{\mathbf{E} \times \mathbf{B}}{B^2}$$



Chappell (1972)

# Energetic particles in the inner magnetosphere

$$\mathbf{u}_{j\perp} = \frac{\mathbf{E} \times \mathbf{B}}{B^2} + \frac{\mathbf{B}}{q_j n B^2} \times \nabla p_j$$

$j = \text{ions, electrons}$

- MHD ordering  
ion thermal speed  $\sim \text{ExB drift} \gg \text{diamagnetic drift}$
- Drift ordering  
ion thermal speed  $\gg \text{ExB drift} \sim \text{diamagnetic drift}$

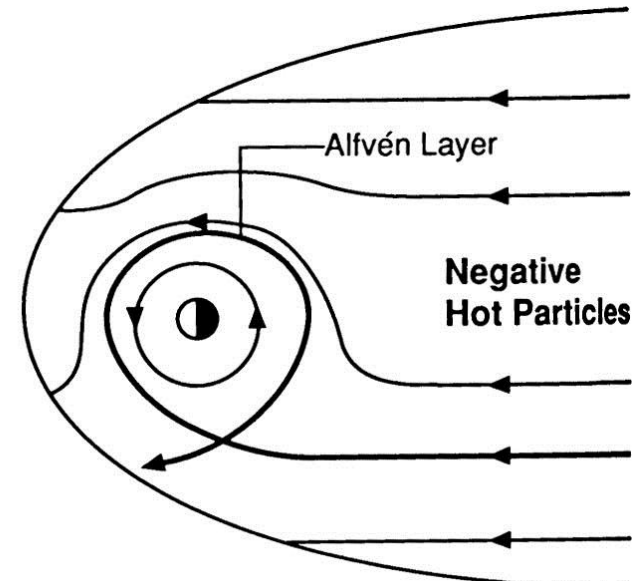
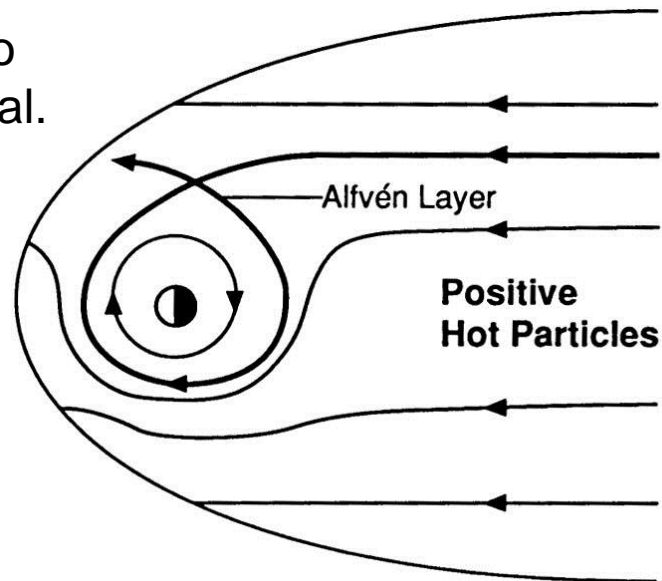
$$\mathbf{u}_{j\perp} = \frac{\mathbf{B} \times \nabla \Psi}{B^2} + \frac{\mathbf{B} \times \nabla \Psi_{\mu j}}{B^2}$$

$$\Psi_{\mu j} = \frac{\mu_j B_0}{q_j L^3}$$

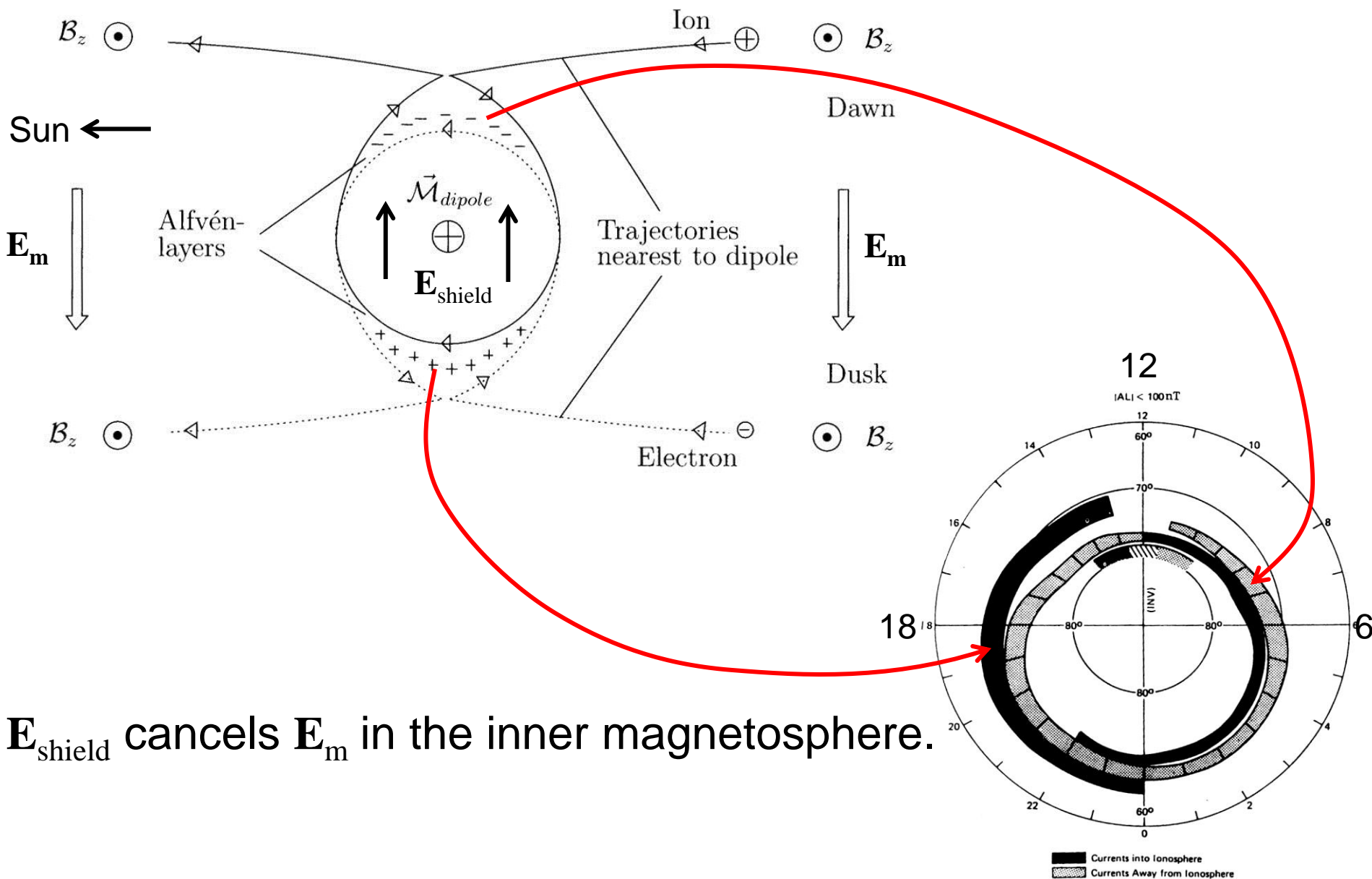
$$\mu_j = \frac{p_j}{nB} = \text{const.}$$

There is a similarity to the corotation potential.

$$\Psi_{cor} = -\frac{a^2 \Omega B_0}{L}$$



# Charge separation in the inner magnetosphere



## Review

- What is the role of the Chapman-Ferraro current?
- What is the magnetotail?
- The idea of “convection” in the magnetosphere-ionosphere system is not simply the plasma bulk flow  $\mathbf{u}$ . How is it different from simple  $\mathbf{u}$ ?
- What is a dynamo?